

ON A PROBLEM OF EQUILIBRIUM OF ELASTIC HALF-SPACE

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The problem regarding the action, on the boundary of a half-space along an infinite straight line, of normal loading varying in accordance with a linear law, is considered.

Let the boundary of elastic half-space be the plane $x = 0$, the x -axis being directed along the internal normal. The loading is distributed along the line $y = 0$ in accordance with the law

$$p_x = p_0 x \tag{1}$$

At the point ($x = \xi$, $y = z = 0$) we isolate from the distributed loading the elementary force $dP_\xi = p_\xi d\xi = p_0 \xi d\xi$. The components of elementary displacement produced by this force, on the basis of Boussinesq's solution ([1], p. 157) are

$$\begin{aligned} du_\xi &= \frac{p_0}{4\pi\mu} \frac{(x-\xi) z \xi d\xi}{R_1^3} - \frac{p_0}{4\pi(\lambda-\mu)} \frac{(x-\xi) \xi d\xi}{R_1(R_1+z)} & (R_1^2 = (x-\xi)^2 + y^2 + z^2) \tag{2} \\ dv_\xi &= \frac{p_0}{4\pi\mu} \frac{y z \xi d\xi}{R_1^3} - \frac{p_0}{4\pi(\lambda+\mu)} \frac{y \xi d\xi}{R_1(R_1+z)}, & dw_\xi = \frac{p_0}{4\pi\mu} \frac{z^2 \xi d\xi}{R_1^3} + \frac{p_0(\lambda+2\mu)}{4\pi\mu(\lambda+\mu)} \frac{\xi d\xi}{R_1} \end{aligned}$$

Integrating expressions (2) in the range of variation of ξ from $-\infty$ to $+\infty$, we obtain the solution of the formulated problem in terms of components of displacements:

$$\begin{aligned} u &= \frac{p_0(\lambda+2\mu)}{2\pi\mu(\lambda+\mu)} z \ln r + \frac{p_0}{2\pi(\lambda+\mu)} y \operatorname{arc} \operatorname{tg} \frac{z}{y} + \frac{p_0 z}{2\pi\mu} & (r^2 = y^2 + z^2) \\ v &= \frac{p_0}{2\pi\mu} \frac{xyz}{r^2} + \frac{p_0}{2\pi(\lambda+\mu)} x \operatorname{arc} \operatorname{tg} \frac{z}{y}, & w = \frac{p_0}{2\pi\mu} \frac{zz^2}{r^2} - \frac{p_0(\lambda+2\mu)}{2\pi\mu(\lambda+\mu)} x \ln r \tag{3} \end{aligned}$$

From here it is easy to obtain the relationship for the components of stress

$$\begin{aligned} X_x &= \frac{2p_0\sigma xz}{\pi r^2} & X_y &= \frac{p_0 yz}{\pi r^2} + \frac{p_0(1-2\sigma)}{\pi} \operatorname{arc} \operatorname{tg} \frac{z}{y}, & Y_y &= -\frac{2p_0 xy^2z}{\pi r^4} \\ Y_z &= -\frac{2p_0 xyz^2}{\pi r^4}, & Z_z &= -\frac{2p_0 xz^3}{\pi r^4} & Z_x &= \frac{p_0 z^2}{\pi r^2} + \frac{p_0(1-\sigma)}{\pi} \end{aligned} \tag{4}$$

Here σ is Poisson's ratio.

The solution of Flamant for the case when the intensity of the distributed loading is equal to p_0 , has the following form in our notation in terms of the components of displacement [1]

$$u^\circ = 0, \quad v^\circ = \frac{p_0}{2\pi\mu} \frac{yz}{r^2} - \frac{p_0}{2\pi(\lambda + \mu)} \operatorname{arc} \operatorname{tg} \frac{y}{z}, \quad w^\circ = -\frac{p_0}{2\pi\mu} \frac{y^2}{r^2} - \frac{p_0(\lambda + 2\mu)}{2\pi\mu(\lambda + \mu)} \ln r \quad (5)$$

By means of elementary transformations these expressions are reduced to the form

$$u^\circ = 0, \quad v^\circ = \frac{p_0}{2\pi\mu} \frac{yz}{r^2} + \frac{p_0}{2\pi(\lambda + \mu)} \operatorname{arc} \operatorname{tg} \frac{z}{y}, \quad w^\circ = \frac{p_0}{2\pi\mu} \frac{z^2}{r^2} - \frac{p_0(\lambda + 2\mu)}{2\pi\mu(\lambda + \mu)} \ln r \quad (6)$$

The following components of stress ([1], p. 207) correspond to these displacements

$$X_x^\circ = \frac{2p_0\sigma}{\pi} \frac{z}{r^2}, \quad Y_y^\circ = -\frac{2p_0}{\pi} \frac{zy^2}{r^4}, \quad Y_z^\circ = -\frac{2p_0}{\pi} \frac{yz^2}{r^4}, \quad Z_z^\circ = -\frac{2p_0}{\pi} \frac{z^3}{r^4} \quad (7)$$

Comparing the obtained solution (3) and (4) with the Flamant solution (6) and (7) we note the following:

$$v = v^\circ x, \quad w = w^\circ x, \quad X_x = X_x^\circ x, \quad Y_y = Y_y^\circ x, \quad Z_z = Z_z^\circ x, \quad Y_z = Y_z^\circ x$$

The analogy between the two problems is not limited by this. In Flamant's problem $u^0 = 0$ indicates the presence of a plane state of strain. In our problem $u \neq 0$, but $du/dx = 0$, the linear deformation is absent along the x -axis, all planes $x = \text{const}$ are equally displaced.

It should be noted that the problem considered may be also be solved by the method of Gutman, suggested especially for problems with a linearly varying loading [2].

BIBLIOGRAPHY

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